# Synthesis and Crystal Structure of $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ 

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#### Abstract

Hydrothermal synthesis starting from $\beta-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ leads to $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ : orthorhombic, space group $P b 2_{1} a ; Z=4 ; a=6.830(1) \AA, b=9.233(2) \AA, c=8.817(2) \AA ; R=0.037$ for 3371 reflections. $\mathrm{NH}_{4}^{+}$cations connect isolated chains of trans-corner sharing $\mathrm{VO}_{5}$ square pyramids running along the $a$-axis. The $\alpha$-phase transforms irreversibly into the $\beta$-form by heating near $200^{\circ} \mathrm{C}$. Comparison is made between the two phases. © 1992 Academic Press, Inc.


## Introduction

The characterization of any phase in the $\mathrm{V}-\mathrm{P}-\mathrm{O}$ system, catalysts as well as their precursors, is of particular interest owing to the catalytic properties of some oxovanadium phosphates in the maleic anhydride obtainment from $\mathrm{C}_{4}$ hydrocarbons (1, 2). Advances in the knowledge of the solid state chemistry of some precursors in the V-P-O-H system were recently reviewed (3, 4); the crystal structures of $\alpha$ - and $\beta-\mathrm{VO}_{2} \mathrm{HPO}_{4} \cdot 2 \mathrm{H}_{2} \mathrm{O}(5,6)$ were determined ab initio from X-ray powder data. Another precursor of the ( VO$)_{2} \mathrm{P}_{2} \mathrm{O}_{7}$ active catalyst is $\beta-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$, as shown by Pulvin et al. $(7,8)$; this compound belongs to a family of salts having the general formulation $A\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)\left(A=\mathrm{NH}_{4}^{+}, \mathrm{K}^{+}, \mathrm{Rb}^{+}\right.$, and $\mathrm{Cs}^{+}$) earlier reported by Preuss and Schug (9). We recently reported the crystal
structure determination of these salts from X-ray powder data (10). The presence of the monohydrogenophosphate group was in fact established at this occasion, since it was previously reported as $A(\mathrm{VOOH})\left(\mathrm{PO}_{4}\right)$ (7) or $\mathrm{AHVPO}_{6}$ (9); its structure (space group Pbca, $Z=8, a=6.8064(6) \AA, b=9.2567(7)$ $\AA$, and $c=17.732(2) \AA$ when $A=\mathrm{NH}_{4}^{+}$) is built up from isolated infinite chains of trans-corner sharing $\mathrm{VO}_{5}$ square pyramids running along the $a$-axis.

We report here on the synthesis and the structure determination, from single crystal X-ray diffraction data, of a new polymorph: $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$.

## Experimental

A $1.3-M$ solution of $\beta-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$, synthesized as described in (9), was evaporated to a smaller volume until incipient
cloudiness. A $15-\mathrm{cm}^{3}$ portion of this suspension was introduced in a hermetic tefion jar placed in a metallic container (PARR bomb) and heated at $200^{\circ} \mathrm{C}$ for 24 hr . The apparatus was allowed to cool down to room temperature for 48 hr . The resulting yellow crystals were filtered out and washed with small portions of acetone.

The transition $\alpha-\beta-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ was characterized by a DSC experiment (DS2 Perkin-Elmer) realized under a flowing $\mathrm{N}_{2}$ atmosphere. A crucible containing 4.30 mg of $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ crystals was heated at $20^{\circ} \mathrm{C} \mathrm{min}^{-1}$ up to $560^{\circ} \mathrm{C}$. The differentiation of the phases (from several DSC stopped at the desired temperature) was realized by powder X-ray diffractometry (Siemens D501). The DSC curve presented two endothermic effects: one in the temperature ranges $186-213^{\circ} \mathrm{C}$ and the other in $294-395^{\circ} \mathrm{C}$. The first peak ( $3.11 \mathrm{cal} \mathrm{g}^{-1}$ ) was attributed to the $\alpha-\beta$ irreversible transition and the second one to the elimination of $\mathrm{NH}_{3}$. After cooling from $560^{\circ} \mathrm{C}$ to room temperature, the material was amorphous. The $\alpha \rightarrow \beta$ transition also can be induced under mechanical pressure ( $\geq 2$ Ton $\mathrm{cm}^{-2}$ ).

## Structure Determination

A small yellow platelet limited by $100 /$ $010 / 001$ faces was selected for the structure determination. The data were collected on a Siemens AED2 four-circle diffractometer. Table I gathers the conditions of the experiment. The cell parameters were refined from 28 reflections well distributed in the reciprocal space at $\approx 30^{\circ} 2 \theta$. The observed conditions limiting possible reflections were consistent with the Pbma and $P b 2_{1} a$ space groups (this non-standard representation was chosen in order to facilitate the comparison with $\beta-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ which is described using the Pbca space group). All the calculations were made with the SHELX-76 program (11); atomic scattering factors and anomalous dispersion correction were taken
from the International Tables for X-Ray Crystallography (12). The $E$ statistic was clearly in favor of the acentric $P b 2_{1} a$ space group $\left(\left\langle E^{2-} 1\right\rangle=0.746\right)$, however, we first tried to solve the structure using the centric Pbma space group. A starting set of atomic coordinates with V and P on $4 d$ sites was obtained from the direct methods. Successive refinements and Fourier difference syntheses did not allow us to complete a fully credible structure model: some interatomic distances and thermal motions were unrealistic, the minimal $R$-factor was 0.17 . Refinements using the $P b 2_{1} a$ space group converged well, starting from the previous set of positions. With all V, P, N, and O atoms, the $R$-factor decreased to 0.084 in the isotropic thermal motion approximation ( $36 \mathrm{pa}-$ rameters refined). Refinements with anisotropic thermal motion lead to the residual $R=0.040$. At this stage, five independent hydrogen atoms on general position were located from $\mathrm{a} \sin \theta_{\text {max }}=0.4$ limited Fourier difference synthesis. Refinements of their coordinates together with a common isotropic $B$-factor lead to the final residuals listed Table I. Refinements after changing all $y$ coordinates into $-y$ lead to $R=0.046$, $R_{w}=0.053$, confirming the absolute configuration choice. $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ really shows an important deviation from a centrosymmetric description in the Pbma space group, since it would have required at least V and P atoms to be on the special position $4 d$ with the same $y=\frac{1}{4}$ coordinate.

The atomic coordinates and thermal parameters are gathered in Table II; selected bond lengths and angles are listed in Table III. The list of structure factors may be obtained on request to the authors.

## Description of the Structure and Discussion

The structure of $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ is built up from $\mathrm{VO}_{5}$ square pyramids linked by trans-vertices through the $\mathrm{O}(1)$ atoms, in

TABLE I
Crystallographic Data for $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$

| Formula weight | 196.96 |
| :---: | :---: |
| Space group | Pb2 ${ }_{1} a$ |
| Z | 4 |
| Calculated density | 2.353 |
| Radiation | MoK $\alpha^{\prime}$ |
| Monochromator | graphite |
| $\mu\left(\mathrm{cm}^{-1}\right), \mathrm{MoK} \alpha$ | 19.7 |
| Crystal size (mm) | $0.120 \times 0.247 \times 0.027$ |
| Lattice constants | $a(\AA)=6.830(1)$ |
|  | $b(\AA)-9.233(2)$ |
|  | $c(\AA)=8.817(2)$ |
|  | $V\left(\AA^{3}\right)=556.01$ |
| Scan type | w- $2 \theta$ |
| $T$ | $20^{\circ} \mathrm{C}$ |
| Angular range ( ${ }^{\circ} 2 \theta$ ) | 4.4-90.0 |
| Learn profile data collection |  |
| Isotropic line width $w=(A+B \operatorname{tg} \theta)^{\circ}$ | $A=0.76$ |
|  | $B=0.33$ |
| Maximum $h, k, l$ | 13, 18, 17 |
| Data examined | 8068 |
| Unique reflections ( $m 2 m$ ) | 4342 |
| $R$ average | 0.031 |
| Merged data retained $I>3 \sigma_{(f)}$ | 3371 |
| Absorption correction | Gauss method |
| Min. and Max. transmission factors | 0.779, 0.932 |
| No extinction correction |  |
| Weighting scheme | $w=1.1 /\left[\sigma^{2}\left(\left\|F_{0}\right\|\right)+0.0020 F_{0}^{2}\right]$ |
| Discrepancy factors | $R=0.037, R_{w}=0.044$ |
| Parameters refined | 97 |
| Max. Shift/esd | 0.001 |
| Max. and Min. electron density in final Fourier difference map | 1.40, $-0.97 \mathrm{e}^{-3}$ |

order to form infinite isolated chains running along the $a$-axis. The $\mathrm{PO}_{4}$ tetrahedra alternate along the chain and link to the remaining trans-vertices of the pyramids (namely $O(2), O(3)$ ). Figure 1 shows that two $\left[\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)\right]_{n}^{n-}$ chains are stacked along the $b$-axis in such a way that the free apices ( $O(5)$ oxygen atoms) of the square pyramids, which are always pointing in the same $b$-direction, would build $\mathrm{VO}_{6}$ octahedron by completion of the square pyramids of the next chain. The resulting $\mathrm{V}-\mathrm{O}(5)$ distance ( $3.110 \AA$ ) is too long to be considered as representative of a normal bond, how-
ever, the topology is consequently near that of isolated perovskite planes parallel to the $a b$-plane. One may consider that $\alpha$ $\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ presents a superstructure since, without the opposite tilting which differentiates two adjacent chains stacked on the $a c$-plane, the $b$-axis would have been divided by a factor 2 . The $\mathrm{NH}_{4}^{+}$ion links the chains by a clear hydrogen bonding scheme. Two hydrogen atoms fix the $\mathrm{NH}_{4}^{+}$ion to one given chain through $\mathrm{H}(4) \cdots \mathrm{O}(1)$ and $\mathrm{H}(1) \cdots \mathrm{O}(6)$ bonds; the two remaining hydrogen atoms are connected to two other different chains: one neighboring chain in

TABLE II
Atomic Coordinates and Thermal Parameters $\left(\mathrm{U}_{\mathrm{ij}} \times 10^{4}\right)^{a}$ For $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$, e.s.d. in Parentheses

| Atom | $x$ | $y$ | $z$ | $U_{11}$ | $U_{22}$ | $U_{33}$ | $\mathrm{U}_{23}$ | $U_{13}$ | $U_{12}$ | $\mathrm{B}_{\mathrm{eq}}\left(\AA^{2}\right)$ or $B$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| V | $0.02856(4)$ | 0 | 0.51851(3) | 91(1) | 191(1) | 120(1) | 8(1) | 1(1) | O(1) | 1.06 (1) |
| P | $0.31578(7)$ | 0.06376 (7) | 0.80033 (5) | 117(2) | 160(2) | 109(1) | -13(2) | 0(1) | -7(2) | 1.02(1) |
| N | $0.7796(4)$ | 0.2582(3) | 0.8059(3) | 257(9) | 263(9) | 219(8) | -39(7) | 33(7) | -51(8) | 1.94(7) |
| $\mathrm{O}(1)$ | 0.2944(2) | $0.0538(2)$ | $0.4470(2)$ | 108(4) | 304(8) | 151(5) | 42(6) | 3(4) | 2(5) | 1.48(4) |
| O(2) | 0.4845 (2) | $0.0085(2)$ | $0.7000(2)$ | 149(5) | 279(7) | 130(4) | -7(7) | 10(4) | 56(6) | 1.47(4) |
| $\mathrm{O}(3)$ | $0.1256(2)$ | 0.0763 (3) | $0.7077(2)$ | 128(5) | 420(9) | 163(5) | -53(7) | -15(4) | 41 (6) | 1.87(5) |
| $\mathrm{O}(4)$ | 0.2842(3) | $0.9486(2)$ | 0.9282(2) | 316(9) | 242(7) | 180(6) | 42(6) | 7(6) | -99(7) | 1.94(6) |
| $\mathrm{O}(5)$ | $0.0320(4)$ | 0.8288(3) | 0.5418(3) | 322(9) | 214(7) | 302(9) | 63(7) | 10(7) | 37(7) | 2.20 (7) |
| $\mathrm{O}(6)$ | $0.3712(3)$ | $0.2026(2)$ | $0.8785(2)$ | 235(7) | 167(6) | 229(7) | -43(5) | -29(6) | 0 (5) | 1.66 (5) |
| H(1) | 0.650(9) | 0.259(6) | 0.808(7) | 379(68) |  |  |  |  |  | $3.0(5)$ |
| H(2) | 0.824(9) | 0.226 (7) | 0.908(7) | 379(68) |  |  |  |  |  | $3.0(5)$ |
| H(3) | 0.836(9) | 0.327 (7) | 0.781 (7) | 379(68) |  |  |  |  |  | $3.0(5)$ |
| H(4) | 0.793(9) | 0.211(7) | 0.722(7) | 379(68) |  |  |  |  |  | $3.0(5)$ |
| H(5) | $0.241(9)$ | 0.866(7) | 0.912(6) | 379(68) |  |  |  |  |  | $3.0(5)$ |

${ }^{a}$ Vibrational $U_{i j}$ coefficients relate to the expression $T=\exp \left[-2 \pi^{2}\left(h^{2} a^{* 2} U_{11}+\cdots+2 h k a^{*} b^{*} U_{12}\right)\right]$.
the $b$-direction through $\mathrm{H}(3) \cdots \mathrm{O}(2)$ and one neighboring chain in the $c$-direction through $\mathrm{H}(2) \cdots \mathrm{O}(6)$. The $\mathrm{H}(5)$ hydrogen atom of the $\mathrm{HPO}_{4}$ group also links two chains by connecting two neighboring $\mathrm{PO}_{4}$ tetrahedra in the $b$-direction through $\mathrm{P}-\mathrm{O}(4)-$ $\mathrm{H}(5) \cdots \mathrm{O}(6)-\mathrm{P}$ bond. The terminal $\mathrm{O}(6)$ oxygen atom of the $\mathrm{HPO}_{4}$ group is consequently involved in three hydrogen bonds.

A valence bond analysis (Table IV) is in agreement with the proposed hydrogen bonding scheme. The analysis was carried out using the Zachariasen law for the V-O and $\mathrm{P}-\mathrm{O}$ bonds $\left(s=\exp \left[\left(R_{\mathrm{o}}-R\right) / B\right]\right.$ with $R_{\mathrm{o}}=1.62, B=0.36$ for P , and $R_{\mathrm{o}}=1.79$, $B=0.319$ for V (13)); owing to the large uncertainty in the observed $\mathrm{N}-\mathrm{H}$ and $\mathrm{O}-\mathrm{H}$ distances, we attributed, somewhat arbitrarily, 0.75 valence units (v.u.) to $\mathrm{N}-\mathrm{H}$ and $\mathrm{O}-\mathrm{H}$ bonds and $0.25 \mathrm{v} . \mathrm{u}$. to $\mathrm{H} \cdots \mathrm{O}$ bonds. From the results presented Table IV, one must conclude that the repartition of the $1^{+}$ charge of $\mathrm{NH}_{4}$ ions through the four $\mathrm{H} \cdots \mathrm{O}$ bonds, each corresponding to 0.25 v.u., is a very imperfect model: in fact, the three different oxygen atoms contributing to $\mathrm{O} \cdots \mathrm{H}-\mathrm{N}$ bonds (namely $\mathrm{O}(1), \mathrm{O}(2), \mathrm{O}(6)$ ) show all an excess of nearly 0.1 v.u.; at variance the $O(3), O(4)$ and $O(5)$ oxygen
atoms, which may be considered as belonging to the $\mathrm{NH}_{4}$ coordination sphere (see distances Table III), show a nearly 0.1 v.u. deficit. Nevertheless, the $V$ and $P$ calculated valences match fairly well the expected values of 5 v.u. in both cases; this allows us to confirm that $\mathrm{V}-\mathrm{O}$ and $\mathrm{P}-\mathrm{O}$ distances are quite normal in this structure.

The occurence of two very short V-O distances ( 1.594 and $1.702 \AA$ ) must be emphasized, although this is not uncommon for $\mathrm{V}^{5+}$ in square pyramid coordination (for instance in $\mathrm{V}_{2} \mathrm{O}_{5}$ (14), the $\mathrm{V}-\mathrm{O}$ distances are $1.577,1.779,1.878 \times 2$, and $2.017 \AA$ ). The shortest V-O distance is characteristic of the double $\mathrm{V}=\mathrm{O}$ bond; the other short distance corresponds to the connection between two $\mathrm{VO}_{5}$ pyramids. Very similar results were observed in the case of $\beta-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)(10)$, in spite of the medium accuracy which characterizes its structure determined and refined from X-ray powder data.

The structural relationship between $\alpha$ and $\beta-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ becomes obvious when Fig. 2 is compared to Fig. 1. One can see that the fundamental building block, defined by the whole $\left[\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)\right]_{n}^{n-}$ infinite chain, is common to the two phases. The

TABLE III
Selected Interatomic Distances $(\AA)$ and Angles $\left({ }^{\circ}\right)$ For $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$, e.s.d. in Parentheses

| V Square pyramid $\langle\mathrm{V}-\mathrm{O}\rangle=1.832$ |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| V | $\mathrm{O}(5)$ | $\mathrm{O}(1)$ | $\mathrm{O}(3)$ | $\mathrm{O}(2)$ | $\mathrm{O}(1)$ |
| $\mathrm{O}(5)$ | 1.594(2) | $2.638(3)$ | 2.788(2) | $2.721(2)$ | 2.868(2) |
| $\mathrm{O}(1)$ | 106.3(2) | 1.702(1) | 2.650(2) | $2.615(2)$ | 3.541(2) |
| $\mathrm{O}(3)$ | 104.2(2) | 93.5(1) | 1.928(1) | $3.774(3)$ | 2.580(2) |
| $\mathrm{O}(2)$ | 99.7(2) | 91.1(2) | 153.2(2) | 1.952(1) | 2.517(2) |
| $\mathrm{O}(1)$ | 106.0(2) | 147.5(1) | 82.5(2) | 79.5(1) | 1.985(1) |
| P Tetrahedron $\langle\mathbf{P}-\mathrm{O}\rangle=1.537$ |  |  |  |  |  |
| P | $\mathrm{O}(6)$ |  | $\mathrm{O}(3)$ | $\bigcirc(2)$ | $\mathrm{O}(4)$ |
| O(6) | 1.504(2) |  | 2.538(2) | 2.507(1) | 2.459(2) |
| $\mathrm{O}(3)$ | 113.1(2) |  | 1.539(1) | 2.531(2) | 2.519(2) |
| $\mathrm{O}(2)$ | 110.9(1) |  | 110.6(2) | 1.540(1) | $2.495(2)$ |
| $\mathrm{O}(4)$ | 106.5(2) |  | 108.5(2) | 107.0(2) | 1.565(1) |
| N Tetrahedron $\langle\mathrm{N}-\mathrm{H}\rangle=0.880$ |  |  |  |  |  |
| N | H(3) |  | H(4) | H(1) | H(2) |
| H(3) | 0.77(6) |  | 1.23(8) | 1.44(8) | 1.46(6) |
| H(4) | 97(9) |  | 0.86(6) | 1.31 (8) | 1.66(9) |
| $\mathrm{H}(1)$ | 120(9) |  | 97(9) | 0.88(6) | 1.51(8) |
| H(2) | 110(9) |  | 126(9) | 107(9) | 0.99(6) |

Hydrogen bonding scheme

| $\mathrm{O} / \mathrm{N}-\mathrm{H} \cdots \mathrm{O}$ | $\mathrm{O} / \mathrm{N}-\mathrm{H}$ | $\mathrm{H} \cdots \mathrm{O}$ | $\mathrm{O} / \mathrm{N}-\mathrm{H} \cdots \mathrm{O}$ | $\mathrm{O} / \mathrm{N}-\mathrm{O}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{O}(4)-\mathrm{H}(5) \cdots \mathrm{O}(6)$ | $0.83(6)$ | $1.72(6)$ | $174(7)$ | $2.545(2)$ |
| $\mathrm{N}-\mathrm{H}(1) \cdots \mathrm{O}(6)$ | $0.88(6)$ | $2.07(6)$ | $157(6)$ | $2.908(3)$ |
| $\mathrm{N}-\mathrm{H}(2) \cdots \mathrm{O}(6)$ | $0.99(6)$ | $2.09(5)$ | $166(6)$ | $2.898(3)$ |
| $\mathrm{N}-\mathrm{H}(3) \cdots \mathrm{O}(2)$ | $0.77(6)$ | $2.20(6)$ | $175(6)$ | $2.968(3)$ |
| $\mathrm{N}-\mathrm{H}(4) \cdots \mathrm{O}(1)$ | $0.86(6)$ | $2.08(6)$ | $165(6)$ | $2.923(3)$ |

$\mathrm{N}-\mathrm{O}$ Distances

| $\mathrm{N}-\mathrm{O}(6)$ | $2.898(3)$ |
| :--- | :--- |
| $\mathrm{N}-\mathrm{O}(6)$ | $2.908(3)$ |
| $\mathrm{N}-\mathrm{O}(1)$ | $2.923(3)$ |
| $\mathrm{N}-\mathrm{O}(4)$ | $2.963(3)$ |
| $\mathrm{N}-\mathrm{O}(2)$ | $2.968(3)$ |
| $\mathrm{N}-\mathrm{O}(3)$ | $3.026(3)$ |
| $\mathrm{N}-\mathrm{O}(2)$ | $3.202(3)$ |
| $\mathrm{N}-\mathrm{O}(5)$ | $3.221(4)$ |
| $\mathrm{N}-\mathrm{O}(5)$ | $3.388(4)$ |



$$
\begin{aligned}
& V \text { at } y=0.00 \\
& P \text { at } y=0.06 \\
& N \text { at } y=0.26
\end{aligned}
$$




$$
\begin{aligned}
& V \text { at } y=0.50 \\
& P \text { at } \dot{y}=0.56 \\
& N \text { at } y=0.76
\end{aligned}
$$

Fig. 1. Projections along $b$ of $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$. The hydrogen bonding scheme is shown on the left part of the figure. The $\mathrm{O}(4)-\mathrm{H}(5) \cdots \mathrm{O}(6)$ are almost superimposed and bonds are not marked. Note that the free apices $(\mathrm{O}(5))$ of the $\mathrm{VO}_{5}$ square pyramids are all situated behind the vanadium atoms. ORTEP drawing (15), ellipsoids scaled to include $90 \%$ probability, thermal motion of H atoms reduced by a factor of 10 .

TABLE IV
Valence Bond Analysis for $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$

|  | V | P | H(1) | H(2) | H(3) | H(4) | H(5) | $\Sigma$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{O}(1)$ | 1.32 |  |  |  |  | 0.25 |  | 2.11 |
|  | 0.54 |  |  |  |  |  |  |  |
| O(2) | 0.60 | 1.25 |  |  | 0.25 |  |  | 2.10 |
| O(3) | 0.65 | 1.25 |  |  |  |  |  | 1.90 |
| O(4) |  | 1.17 |  |  |  |  | 0.75 | 1.92 |
| O(5) | 1.85 |  |  |  |  |  |  | 1.85 |
| O(6) |  | 1.38 | 0.25 | 0.25 |  |  | 0.25 | 2.13 |
| N |  |  | 0.75 | 0.75 | 0.75 | 0.75 |  | 3.00 |
| $\Sigma$ | 4.96 | 5.05 | 1.00 | 1.00 | 1.00 | 1.00 | 1.00 |  |



V at $y= \pm 0.54$
$P$ at $y= \pm 0.60$
N at $\mathrm{y}= \pm 0.19$

$V$ at $y= \pm 0.04$
$P$ at $y= \pm 0.10$
N at $\mathrm{y}= \pm 0.31$

Fig. 2. Projections along $b$ of $\beta-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$. Note that the free apices of the $\mathrm{VO}_{5}$ square pyramids are situated behind the vanadium atoms for $0<z<\frac{1}{2}$ and the contrary for $\frac{1}{2}<z<1$. ORTEP drawing (15).
stacking of the chains along the $b$-axis is also identical: the main difference is that the $c$-axis is doubled in the $\beta$-phase where the chains at $z \approx \frac{1}{4}$ show significantly different $y$-coordinates than those at $z \approx \frac{3}{4}$ (the shift is $\Delta y \approx 0.1$ ); moreover, the apex of the $\mathrm{VO}_{5}$ pyramids at $z \approx \frac{1}{4}$ and $z \approx \frac{3}{4}$ point toward opposite directions along the $b$-axis due to the presence of an inversion center. Hydrogen atoms could not be located in the $\beta$-phase, but from the fact that the cell of the $\beta$-phase is built up from two cells of the $\alpha$-phase stacked along the $c$-axis and related by an inversion center, it can be proposed that the hydrogen bonding scheme of the two phases is the same with the exception of the unique hydrogen bond along the $c$-axis (H(2) $\cdots \mathrm{O}(6)$ in the $\alpha$-phase). Such similitudes suggest that the $\alpha \rightarrow \beta$ transition
mechanism is topotactic. In this hypothesis, the reorganization could be realized by a shear, parallel to the $a b$-plane, of $\alpha$-blocks of infinite size along $a_{\alpha}$ and $b_{\alpha}$ and of thickness equal to the $c_{\alpha}$-axis. The displacement of each block is quite small: approximately $0.4 \AA$ in opposite direction along the $b_{\alpha}$-axis; however, in half such blocks, the double bond $\mathrm{V}=\mathrm{O}(5)$ must be formally broken and reinstalled in the opposite direction along the $b_{\alpha}$-axis. This operation consists in a weak displacement of $O(5)$ and a shift of the vanadium atom from one side of the equatorial plane to the other. Crystallographically, there is no objection to such a topotactic transitional model: the $P b 2_{1} a$ space group of the $\alpha$-phase belongs to the maximal non-isomorphic subgroups of the Pbca space group of the $\beta$-phase. The main


Fig. 3. Cell parameters vs ionic radii plot for the $\beta-\mathrm{A}\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)\left(A^{+}=\mathrm{NH}_{4}^{+}, \mathrm{K}^{+}, \mathrm{Rb}^{+}, \mathrm{Cs}^{+}\right)$compounds ( $r$ : ionic radii for sixfold coordination).
difference between the $\alpha_{\mathrm{I}^{-}}$and $\alpha_{\mathrm{II}^{-}}-\mathrm{VOPO}_{4}$ forms (16) is due to such shifts of the vanadium atoms, however, no phase transition has been pointed out.

## Conclusion

The non-existence of either sodium and lithium compounds with the stoichiometry $A\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)\left(A^{+}=\mathrm{Na}^{+}, \mathrm{Li}^{+}\right)$or alkaline compounds showing the $\alpha-\mathrm{NH}_{4}\left(\mathrm{VO}_{2}\right)$ $\left(\mathrm{HPO}_{4}\right)$ structure-type can be now explained in the light of the structural analysis of the two ammonium polymorphs.

Only the $\mathrm{NH}_{4}^{+}$cation shows the two polymorphic ( $\alpha$ and $\beta$ ) varieties of the $A\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$ salts. The structure of both
polymorphs is well described as formed by polymerization of cis- $\mathrm{VO}_{2}^{+}$units by $\mu$-(O, $\left.\mathrm{O}^{\prime}\right)-\left[\mathrm{HOPO}_{3}\right]$ groups. The differences are due to the different packing between these chains and the $\mathrm{NH}_{4}^{+}$cations. It seems that the capability of the $\mathrm{NH}_{4}^{+}$ions to form such a similar network of directed bonds through its hydrogen atoms allows the existence of both $\alpha$ - and $\beta$-phases. This indirectly would justify the failure in the preparation of the $\alpha$-phase in the case of such spherical cations as $\mathrm{K}^{+}, \mathrm{Rb}^{+}$, and $\mathrm{Cs}^{+}$.

Electrostatic interactions involving nonspheric cations such as the ammonium are less effective in order to give close-packed lattices. In fact, the cell volume of $\beta$-polymorphs with spherical cations of similar ionic radius as $\mathrm{K}^{+}$or even larger than $\mathrm{NH}_{4}^{+}$as $\mathrm{Rb}^{+}$are significantly smaller than the cell volume measured for $\beta-\mathrm{NH}_{4}$ $\left(\mathrm{VO}_{2}\right)\left(\mathrm{HPO}_{4}\right)$.

On the other hand, failure in the prepara-


Fig. 4. Variation of the tilting angle in the $\beta$-phases vs the ionic radii.
tion of both $\beta$ - and $\alpha$-phases of lighter alkaline metals can be easily explained from the structural data. Thus from Figs. 3 and 4, a good correlation is found between cell parameters or tilting angle (observed between the medium planes defined by the four basal oxygen atoms of the $\mathrm{VO}_{5}$ pyramids) and the ionic radii (17) of alkaline cations in the $\beta$-phases. Extrapolation down smaller cations such as $\mathrm{Na}^{+}$would imply very strong repulsion among chains resulting in a compression of approximately 0.30 and $0.75 \AA$ in the [010] and [001] directions, respectively, and very distorted topology of the bridging $\mathrm{HPO}_{4}$ groups due to the higher expected value of the tilting angle ( $26.6^{\circ}$ ).

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